

# 25

## Complex numbers

### Answers to additional problems

**25.1**  $\mathbf{I}_x \mathbf{I}_y = \frac{1}{2} \times \frac{1}{2i} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \frac{1}{4i} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$

$$\mathbf{I}_y \mathbf{I}_x = \frac{1}{2i} \times \frac{1}{2} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \frac{1}{4i} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\text{so } \mathbf{I}_x \mathbf{I}_y - \mathbf{I}_y \mathbf{I}_x = \frac{1}{4i} \begin{pmatrix} -2 & 0 \\ 0 & 2 \end{pmatrix} = -\frac{2}{4i} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

The factor of 2/4 cancels to  $\frac{1}{2}$ . We then manipulate to remove the  $i$  from the denominator, saying  $i/i = 1$ ,

$$\mathbf{I}_x \mathbf{I}_y - \mathbf{I}_y \mathbf{I}_x = -\frac{i}{2i^2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = +\frac{i}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = i \times \mathbf{I}_z$$

**25.2**  $\mathbf{I}_y \mathbf{I}_z = \frac{1}{4i} \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} \quad \mathbf{I}_z \mathbf{I}_y = \frac{1}{4i} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\text{so } \mathbf{I}_y \mathbf{I}_z - \mathbf{I}_z \mathbf{I}_y = \frac{1}{4i} \left\{ \begin{pmatrix} 0 & -1 \\ -1 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right\} = \frac{1}{4i} \begin{pmatrix} 0 & -2 \\ -2 & 0 \end{pmatrix} = -\frac{2}{4i} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Again, we cancel the 2 and 4 to yield  $\frac{1}{2}$ . Manipulation to remove  $i$  from the denominator gives,

$$\mathbf{I}_y \mathbf{I}_z - \mathbf{I}_z \mathbf{I}_y = -\frac{i}{2i^2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \frac{i}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = i \times \mathbf{I}_x$$

**25.3**  $\mathbf{I}_z \mathbf{I}_x - \mathbf{I}_x \mathbf{I}_z = \frac{1}{4} \left\{ \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} - \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \right\} = \frac{1}{4} \begin{pmatrix} 0 & 2 \\ -2 & 0 \end{pmatrix}$

$$\text{Factorizing yields } = \frac{2}{4} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

Again, we can cancel the 2 and the 4 to yield  $\frac{1}{2}$ .

The matrix should remind us of the matrix in the expression for  $\mathbf{I}_y$ , but with a slightly different factor. But if we multiply by  $i/i = 1$ , we generate  $i$  times  $\mathbf{I}_y$ ,

$$\frac{i}{2i} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i \times \frac{1}{2i} \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i \times \mathbf{I}_y$$

**25.4**  $\frac{1}{Z_{\text{total}}} = \frac{1}{Z_c} + \frac{1}{Z_{R_1}} + \frac{1}{Z_{R_2}} = \frac{1}{1/i\omega C} + \frac{1}{R_1} + \frac{1}{R_2}$

$$\frac{1}{Z_{\text{total}}} = \frac{R_1 R_2 i \omega C + (R_1 + R_2)}{R_1 R_2}$$

so  $Z_{\text{total}} = \frac{R_1 R_2}{R_1 R_2 i \omega C + (R_1 + R_2)}$

**25.5**  $12 - 6i$ .

**25.6** A general square ( $x^2 + y^2$ ) has roots  $(x + yi)$   $(x - yi)$ . Therefore,  $(3a + 7bi)$   $(3a - 7bi)$ .

**25.7** Using the formula in eqn. (25.11),

$$e^{ikx} = \cos kx + i \sin kx \quad \text{and}$$

$$e^{-ikx} = \cos(-kx) + i \sin(-kx) = \cos(kx) - i \sin(kx)$$

Therefore,  $\psi = Ae^{ikx} + Be^{-ikx}$  can be rewritten as,

$$\psi = A(\cos kx + i \sin kx) + B(\cos(kx) - i \sin(kx))$$

$$\psi = (A + B)\cos kx + (A - B)i \sin kx$$

**25.8** Multiply together the two wavefunctions  $\Psi$  and  $\Psi^*$ , so multiply the brackets in the following problem,

$$\Psi\Psi^* = (f + ig)(f - ig)$$

We simplify this problem by recognizing how the two brackets resemble the factors of the difference of two squares.

Therefore, the product  $\Psi\Psi^* = (f^2 - i^2g^2) = (f^2 + g^2)$ .

**25.9** Classically, we know that the kinetic energy  $E_{\text{KE}}$  is,

$$E_{\text{KE}} = \frac{1}{2}mv^2$$

$$E_{\text{KE}} = \frac{mv^2}{2} \times \frac{m}{m} = \frac{(mv)^2}{2m} = \frac{p^2}{2m}$$

Therefore, substituting in for the w momentum operator, we can derive the kinetic energy operator,

$$E_{\text{KE}} = \frac{p^2}{2m}$$

$$E_{\text{KE}} = \frac{-i\hbar \frac{d}{dx} \times -i\hbar \frac{d}{dx}}{2m} \quad (\text{notice the way that, } -i \times -i = -1)$$

$$E_{\text{KE}} = \frac{(-1)^2(i)^2 \hbar^2 \frac{d}{dx} \left( \frac{d}{dx} \right)}{2m}$$

$$\text{so } E_{\text{KE}} = \frac{-\hbar^2}{2m} \frac{d^2}{dx^2}$$

This operator appears in the one-dimensional Schrödinger equation,

$$-\frac{\hbar^2}{2m} \frac{d^2\Psi}{dx^2} + V\Psi = E\Psi$$

**25.10** We know from eqn. (25.13) that,

$$e^{ikx} = (\cos kx + i \sin kx)$$

$$\text{and } e^{-ikx} = (\cos(-kx) + i \sin(-kx)) = (\cos kx - i \sin kx).$$

Therefore, we can rewrite the wavefunction as,

$$\psi = A(e^{ikx} \pm e^{-ikx})$$

$$\psi = A((\cos kx + i \sin kx) \pm (\cos kx - i \sin kx))$$

$$\psi = A((\cos kx \pm \cos kx) + i(\sin kx \mp \sin kx))$$

There are two possible solutions,

- If we take the top line of the  $\pm$  symbols, then  $\psi = 2A \cos kx$ .
- If we take the lower line of the  $\pm$  symbols, then  $\psi = 2iA \sin kx$ .
- The next step is to see which of these forms of the wavefunction best satisfies the boundary conditions.